

# Positive Lyapunov exponent for ergodic Schrödinger Operators

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Let  $\ell^2(\mathbb{Z})$  be the Hilbert space of square summable sequences.  
Let the **potential**  $V : \mathbb{Z} \rightarrow \mathbb{R}$  be a bounded sequence.  
Define the **Schrödinger operator**  $H$  by

$$H : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$$
$$(Hu)(n) = \underbrace{u(n+1) + u(n-1)}_{(\Delta u)(n)} + V(n)u(n). \quad (1)$$

$H$  is a bounded and self-adjoint operator.  
If  $u$  solves  $Hu = Eu$  for some **energy**  $E$ , then

$$\begin{pmatrix} u(n+1) \\ u(n) \end{pmatrix} = A(E, n) \cdot \begin{pmatrix} u(1) \\ u(0) \end{pmatrix} \quad (2)$$

with

$$A(E, n) = \prod_{j=n}^1 \begin{pmatrix} V(j) - E & -1 \\ 1 & 0 \end{pmatrix}. \quad (3)$$

We call  $A(E, n)$  the **n-step transfer matrix**.

Let  $A(E, n)$  be the  $n$ -step transfer matrix

$$A(E, n) = \prod_{j=n}^1 \begin{pmatrix} V(j) - E & -1 \\ 1 & 0 \end{pmatrix}. \quad (4)$$

The Lyapunov exponent  $L(E)$  describes the exponential growth of the norm of  $A(E, n)$ :

$$L(E) = \limsup_{n \rightarrow \infty} \frac{1}{n} \log \|A(E, n)\|. \quad (5)$$

For fixed energy  $E$ :  $0 \leq L(E) \leq \text{const.}$

By the Combes–Thomas estimate, one also has

$$L(E) \gtrsim \text{dist}(E, \sigma(H)). \quad (6)$$

We will be interested in positive Lyapunov exponent  $L(E) > 0$ , on the spectrum  $\sigma(H)$ . In the ergodic case

$$\sigma_{\text{ac}}(H) = \overline{\{E : L(E) = 0\}}^{\text{ess}}. \quad (7)$$

Thus one can think of positive Lyapunov exponent as a criterion for localization.

Instead of looking at general potentials  $V(n)$ , we will restrict our attention to **ergodic potentials**.

So let  $(\Omega, \mu)$  be a probability space.  $T : \Omega \rightarrow \Omega$  an invertible **ergodic** transformation. Ergodic means that if  $g$  satisfies

$$g \circ T = g \tag{8}$$

then  $g$  is almost surely constant. Ergodic implies that for  $A \subseteq \Omega$

$$\frac{1}{n} \#\{1 \leq k \leq n : T^k \omega \in A\} \rightarrow \mu(A), \quad a.e. \omega. \tag{9}$$

Let  $f : \Omega \rightarrow \mathbb{R}$  a bounded function. We define for  $\omega \in \Omega$  and the **coupling**  $\lambda > 0$ :

$$V_{\omega, \lambda}(n) = \lambda f(T^n \omega). \tag{10}$$

From the definition of  $L_{\omega}(E, \lambda)$ , it follows that

$$L_{T\omega}(E, \lambda) = L_{\omega}(E, \lambda). \tag{11}$$

Thus by ergodicity, the value is **almost surely  $\omega$  independent**, and we denote it by  $L(E, \lambda)$ .

## A 'silly' example: two periodic Schrödinger operator

Take  $\Omega = \mathbb{Z}_2 = \{0, 1\}$  with uniform measure.

$$Tx = x + 1 \pmod{2}, \quad (12)$$

and  $f(x) = 2(x - \frac{1}{2})$ . Then the potential is just

$$\dots, \lambda, -\lambda, \lambda, -\lambda, \dots \quad (13)$$

and we have for the **spectrum**

$$\sigma(H_\lambda) = \sigma_{\text{ac}}(H_\lambda) = [-\sqrt{4 + \lambda^2}, -\lambda] \cup [\lambda, \sqrt{4 + \lambda^2}]. \quad (14)$$

And the Lyapunov exponent **vanishes** on it.

## Known Results: 1/2.

- ▶ **Anderson Model**:  $V(n)$  are independent identically distributed random variables.

By **Furstenberg**, **Kotani**:  $L(E, \lambda) > 0$  for all  $E$  and all  $\lambda > 0$ .

By **Pastur–Figotin**:

$$0 < \lambda \ll 1 : \quad L(E, \lambda) \gtrsim \lambda^2. \quad (15)$$

- ▶ **Rotations on the torus**: ( $\alpha$  irrational)

$$\Omega = \mathbb{R}/\mathbb{Z} \cong [0, 1], \quad T_x = x + \alpha \pmod{1}. \quad (16)$$

$V(n) = \lambda f(n\alpha)$  quasi-periodic,  $f$  analytic,

$$\lambda \gg 1 : \quad L(E, \lambda) \gtrsim \log \lambda \quad (17)$$

for all  $E$ . **Herman**, **Sorets–Spencer**, **Bourgain–Goldstein**, **Goldstein–Schlag**, **Bourgain**.

The Lyapunov exponent **vanishes** on the spectrum for **small**  $\lambda$ .  
**Eliasson**, **Bourgain–Jitomirskaya**.

## Known Results: 2/2.

- ▶ **Skew-Shift** ( $\alpha$  irrational)

$$T : [0, 1]^2 \rightarrow [0, 1]^2, \quad T(x, y) = (x + \alpha, x + y) \pmod{1}.$$

$V(n) = \lambda f(n^2 \alpha)$ ,  $f$  analytic. **Bourgain–Goldstein–Schlag**:

$$\lambda \gg 1 : \quad L(E, \lambda) \gtrsim \log(\lambda). \quad (18)$$

**Conjecture**:  $L(E, \lambda) > 0$  for all  $\lambda > 0$ .

Partial results by **Bourgain** for  $0 < \lambda \ll 1$ .

- ▶ **Doubling**

$$T : [0, 1] \rightarrow [0, 1], \quad Tx = 2x \pmod{1}. \quad (19)$$

$V(n) = \lambda f(2^n x)$ . **Chulaevsky–Spencer**:

$$0 < \lambda \ll 1 : \quad L(\lambda, E) \gtrsim \lambda^2. \quad (20)$$

Also conjectured for all  $\lambda > 0$ .

Positivity is known by **Damanik–Killip**, but proof is non-quantitative (they use **Kotani**).

# Main result: Pseudoformulation

## Theorem

Assume an *initial condition at an initial scale  $n$*  then  $L(E) > 0$  for *most energies  $E$* .

I will give a more specific statement later.

Proof by metaproof.

Multiscale analysis:

Use *energy elimination* to pass from *scale  $n$*  to *scale  $100^j n$*  on step  $j$ . □

Now: What kind of initial condition do we want?

Later: Exact statement and some ideas of the proof.

## Choosing an initial condition: First try

Recall

$$L(E) = \limsup_{n \rightarrow \infty} \frac{1}{n} \log \|A(\omega, n, E)\| \quad (21)$$

for almost every  $\omega$ .

First guess of an initial condition:

$$\mu(\{\omega : \frac{1}{n} \log \|A(\omega, n, E)\| < \gamma\}) \leq \frac{1}{4} \quad (22)$$

for  $\gamma > 0$  and  $n$  suitable.

This condition is **necessary** to obtain  $L(E) \gtrsim \gamma$ .

However, it might not be **sufficient**.

Let  $\gamma > 0$  and  $n$  large, assume

$$\mu(\{\omega : \frac{1}{n} \log \|A(\omega, n, E)\| < \gamma\}) \leq \frac{1}{4}. \quad (23)$$

We know

$$A(\omega, 2n, E) = A(T^n \omega, n, E)A(\omega, n, E). \quad (24)$$

Want to use this to go from scale  $n$  to scale  $2n$ . However,

- ▶ We would lose probability.
- ▶ The matrix norm is not multiplicative, but only **submultiplicative**.

**Solution 1:** By **Bourgain–Goldstein–Schlag**:

Use **Large Deviation Estimates** and the **Avalanche Principle**.

**Problem:** Large Deviation Estimates are not proved for doubling at large coupling.

**Comment:** Single energy method.

## Solution 2: Multiscale Analysis

- ▶ Reformulate the initial condition as a condition for **finite volume Green's function**. For  $x, y \in \Lambda \subseteq \mathbb{Z}$  define

$$G_{\omega, \Lambda}(E, x, y) = \langle \delta_x, (H_{\omega, \Lambda} - E)^{-1} \delta_y \rangle \quad (25)$$

where  $H_{\omega, \Lambda}$  is the restriction of  $H_\omega$  to  $\ell^2(\Lambda)$ .

$$|G_{\omega, [0, 2n]}(E, 0, n)| \leq e^{-\gamma n} \iff \frac{1}{n} \log \|A(\omega, n, E)\| \geq \gamma. \quad (26)$$

- ▶ Use the **resolvent equation**  $x \in [a, b] \subseteq \Lambda = [0, N]$

$$\begin{aligned} G_{\omega, \Lambda}(E, 0, x) &= -G_{\omega, \Lambda}(E, 0, a-1)G_{\omega, [a, b]}(E, a, x) \\ &\quad - G_{\omega, \Lambda}(E, 0, b+1)G_{\omega, [a, b]}(E, b, x) \end{aligned}$$

to pass from scale to scale.

- ▶ Use a **Wegner estimate** to avoid losing probability.

Adapted by **Bourgain** to prove positive Lyapunov exponent **without any a priori Wegner estimate**.

This approach has been heavily used for Random Schrödinger operators to prove Anderson Localization.

## The main result

Theorem (arXiv:0905.1791)

Assume the *initial condition* for  $x \in \{1, 2n - 1\}$

$$\mu \left( \left\{ \omega : \sup_{E \in [E_0, E_1]} |G_{\omega, [1, 2n-1]}(E, n, x)| > e^{-\gamma n} \right\} \right) \leq \frac{1}{8}, \quad (27)$$

for  $\gamma \cdot n \gg 1$  (dependent on  $E_1 - E_0$ ). Then there exists a set  $\mathcal{E}_b$  such that

$$|\mathcal{E}_b| \lesssim e^{-c\gamma n} \quad (28)$$

and for  $E \in [E_0, E_1] \setminus \mathcal{E}_b$

$$L(E) \geq \frac{\gamma}{2}. \quad (29)$$

We will now discuss how to prove an initial condition for ergodic Schrödinger operators at large coupling.  
Then sketch some ideas for the proof.

## Proof of the initial condition

Assume for some  $\alpha > 0$  and all  $E \in \mathbb{R}$

$$\mu(\{\omega : |f(\omega) - E| \leq \varepsilon\}) \lesssim \varepsilon^\alpha. \quad (30)$$

For example [analytic](#)  $f : [0, 1] \rightarrow \mathbb{R}$  and  $\mu$  the Lebesgue measure satisfy this assumption. [Łojasiewicz](#).

This implies for any  $N \geq 1$  that for  $\lambda \gg 1$  ( $V_{\omega,\lambda}(n) = \lambda f(T^n \omega)$ )

$$\mu(\{\omega : \text{dist}(E, \{V_{\omega,\lambda}(n)\}_{n=1}^N) < \lambda^{\alpha/2}\}) \leq \frac{1}{16}. \quad (31)$$

By  $\|\Delta\| \leq 2$  and  $H_\omega = \Delta + V_{\omega,\lambda}$ , we thus obtain

$$\mu(\{\omega : \text{dist}(E, \sigma(H_{\omega,\lambda,[1,N]})) < \frac{1}{2}\lambda^{\alpha/2}\}) \leq \frac{1}{16}. \quad (32)$$

Use the [Combes–Thomas estimate](#) to convert this into a Green's function estimate and to obtain the initial condition.

# The doubling map

In particular:

## Corollary

Let  $f : [0, 1] \rightarrow \mathbb{R}$  be analytic.

Let  $\lambda$  be large enough.

Introduce the potential (almost every  $x$ )

$$V(n) = \lambda f(2^n x \pmod{1}). \quad (33)$$

Then

$$|\{E : L(E, \lambda) \leq \frac{1}{5} \log(\lambda)\}| \lesssim e^{-c\lambda^\alpha} \quad (34)$$

for some  $c > 0, \alpha > 0$ .

Hence  $L(E) \geq \frac{1}{5} \log \lambda$  for most energies  $E$ .

Note  $[-2, 2] + f(0) \subseteq \sigma(\Delta + V)$ .

Note the result holds also for example for any toral automorphism  $[0, 1]^2 \rightarrow [0, 1]^2$ .

## Proof: 1 / 5

Recall for  $x, y \in \Lambda \subseteq \mathbb{Z}$

$$G_{\omega, \Lambda}(E, x, y) = \langle \delta_x, (H_{\omega, \Lambda} - E)^{-1} \delta_y \rangle \quad (35)$$

We assume the **initial condition**

$$\mu \left( \left\{ \omega : \sup_{E \in [E_0, E_1]} \sup_{x \in \{1, 2n-1\}} |G_{\omega, [1, 2n-1]}(E, x, n)| > e^{-\gamma \cdot n} \right\} \right) \leq \frac{1}{4}.$$

I will illustrate, how to do multiscale analysis **without a Wegner estimate**. To simplify things, I will assume the above holds for **all**  $\omega$ .

I will go from **scale**  $n$  to **scale**  $100n$ .

## Proof: 2 / 5

Assume for all  $\omega, \gamma \cdot n \gg 1$  and  $x \in \{1, 2n - 1\}$

$$\sup_{E \in [E_0, E_1]} |G_{\omega, [1, 2n-1]}(E, x, n)| \leq \frac{1}{2} e^{-\gamma \cdot n}. \quad (36)$$

Then for  $E \in [E_0, E_1]$

$$\begin{aligned} & G_{\omega, [1, 200n-1]}(E, 1, 100n) \\ &= -G_{\omega, [1, 200n-1]}(E, 1, 99n) \\ &\quad \cdot G_{\omega, [99n+1, 101n-1]}(E, 99n+1, Ln) + \dots \\ &= -G_{\omega, [1, 200n-1]}(E, 1, 99n) \\ &\quad \cdot \underbrace{G_{T^{100n}\omega, [(1, 2n-1)]}(E, 1, n)}_{|\dots| \leq \frac{1}{2} e^{-\gamma \cdot n}} + \dots \end{aligned}$$

Now [iterate!](#)

## Proof: 3 / 5

Assume for all  $\omega$ ,  $\gamma \cdot n \gg 1$  and  $x \in \{1, 2n - 1\}$

$$\sup_{E \in [E_0, E_1]} |G_{\omega, [1, 2n-1]}(E, x, n)| \leq \frac{1}{2} e^{-\gamma \cdot n}. \quad (37)$$

Iterating the above, we obtain

$$|G_{\omega, [1, 200n-1]}(E, 1, 100n)| \leq e^{-100\gamma \cdot n} \cdot \sup_{x, y \in [1, 200n-1]} |G_{\omega, [1, 200n-1]}(E, x, y)|.$$

Since  $H_{\omega, [1, 200n-1]}$  is [self-adjoint](#)

$$\sup_{x, y \in [1, 200n-1]} |G_{\omega, [1, 200n-1]}(E, x, y)| \leq \frac{1}{\text{dist}(E, \sigma(H_{\omega, [1, 200n-1]}))}. \quad (38)$$

Now  $\sigma(H_{\omega, [1, 200n-1]})$  consists of [only  \$\approx 200n\$  elements](#).

## Proof: 4 / 5

We have

$$|G_{\omega,[1,200n-1]}(E, 1, 100n)| \leq \frac{e^{-100\gamma \cdot n}}{\text{dist}(E, \sigma(H_{\omega,[1,200n-1]}))}. \quad (39)$$

with  $\sigma(H_{\omega,[1,200n-1]})$  consisting of **less than  $200n$  elements**.

To obtain

$$|G_{\omega,[1,200n-1]}(E, 1, 100n)| \leq e^{-99\gamma \cdot n}, \quad (40)$$

we only need

$$\text{dist}(\sigma(H_{\omega,[1,200n-1]}), E) \geq e^{-\gamma \cdot n}. \quad (41)$$

Hence, we only need to eliminate a set of  $E$  of measure at most

$$400n \cdot e^{-\gamma \cdot n} \quad (42)$$

for each fixed  $\omega$ .

## Proof: 5 / 5

Let

$$\mathcal{B} = \{(E, \omega) \in [E_0, E_1] \times \Omega : |G_{\omega, [1, 200n-1]}(E, 1, 100n)| > e^{-99\gamma \cdot n}\}.$$

We have

$$(Leb \times \mu)(\mathcal{B}) \leq 400n \cdot e^{-\gamma \cdot n}. \quad (43)$$

Let  $\mathcal{E}_b$  be the set such that

$$\mu \left( \left\{ \omega : |G_{\omega, [1, 200n-1]}(E, 1, 100n)| > e^{-99\gamma \cdot n} \right\} \right) > \frac{1}{4}$$

holds. By [Markov's inequality](#) we can thus conclude that

$$|\mathcal{E}_b| \leq 1600n \cdot e^{-\gamma \cdot n}. \quad (44)$$

We see here why we need  $E_1 - E_0$  to be big compared to  $e^{-\gamma \cdot n}$ .

New  $\gamma$  is  $\frac{99}{100}\gamma$ .

## Further aspects of the proof

We have seen now how to go from all  $\omega$  and all  $E$  to most  $\omega$  and most  $E$ .

We will need to do better:

- ▶ Need to ensure that  $\mathcal{E}_b$  consists of finitely many intervals, so the above procedure can be iterated.  
**Bourgain**: Semialgebraic sets requires analytic functions.  
Here: Counting argument.
- ▶ Need to deal with not for all  $\omega$ .  
This uses ergodicity and an improved iteration of the resolvent equation.

## Further extensions

- ▶ The methods allow me to say something about  $V(n) = \lambda \cos(n^k)$  for all small  $\lambda$ , letting  $k$  be large depending on  $\lambda$ .  
Here I use **randomness** to obtain an **initial condition**, following the methods **Bourgain–Schlag** used for the doubling map at small coupling.
- ▶ Assuming a (weak) **Wegner estimate**, one can remove the exceptional set of energies.  
However, this is an extra assumption.  
See **Gan-K** (arXiv:0906.3300).
- ▶ **Open Problem**: Prove an initial condition for  $V(n) = \lambda \cos(n^2)$  for small  $\lambda > 0$ .