

Math 211

Lecture #24

November 21, 2000

Forced Harmonic Motion

Assume an oscillatory forcing term:

$$y'' + 2cy' + \omega_0^2 y = A \cos \omega t$$

- A is the forcing amplitude
- ω is the forcing frequency

Forced Undamped Motion

$$y'' + \omega_0^2 y = A \cos \omega t$$

Homogeneous equation

$$y'' + \omega_0^2 y = 0$$

General solution

$$y(t) = C_1 \cos \omega_0 t + C_2 \sin \omega_0 t.$$

If $\omega = \omega_0$ we have an exceptional case.

$$\omega \neq \omega_0$$

$$y'' + \omega_0^2 y = A \cos \omega t$$

Look for a particular solution of the form

$$x_p(t) = a \cos \omega t + b \sin \omega t.$$

We find

$$x_p(t) = \frac{A}{\omega_0^2 - \omega^2} \cos \omega t.$$

General solution

$$x(t) = C_1 \cos \omega_0 t + C_2 \sin \omega_0 t + \frac{A}{\omega_0^2 - \omega^2} \cos \omega t.$$

Initial conditions $x(0) = x'(0) = 0 \Rightarrow$

$$x(t) = \frac{A}{\omega_0^2 - \omega^2} [\cos \omega t - \cos \omega_0 t].$$

Example: $\omega_0 = 9$, $\omega = 8$, $A = \omega_0^2 - \omega^2 = 17$.

$$x(t) = \cos 9t - \cos 8t.$$

Set

$$\bar{\omega} = \frac{\omega_0 + \omega}{2} \quad \text{and} \quad \delta = \frac{\omega_0 - \omega}{2}.$$

Then

$$\omega = \bar{\omega} - \delta \quad \text{and} \quad \omega_0 = \bar{\omega} + \delta,$$

and

$$\begin{aligned} x(t) &= \frac{A}{\omega_0^2 - \omega^2} [\cos \omega t - \cos \omega_0 t] \\ &= \frac{A \sin \delta t}{2\bar{\omega}\delta} \sin \bar{\omega} t. \end{aligned}$$

$$x(t) = \frac{A \sin \delta t}{2\bar{\omega}\delta} \sin \bar{\omega}t$$

Example:

$$\bar{\omega} = 8.5 \quad \text{and} \quad \delta = 0.5.$$

- Envelope: Slow oscillation with frequency δ .
- Fast oscillation with frequency $\bar{\omega}$ and varying amplitude.
- Beats.

$$\omega = \omega_0$$

$$y'' + \omega_0^2 y = A \cos \omega_0 t$$

We have an exceptional case. Try

$$x_p(t) = t[a \cos \omega t + b \sin \omega t].$$

We find

$$x_p(t) = \frac{A}{2\omega_0} t \sin \omega_0 t.$$

General solution

$$x(t) = C_1 \cos \omega_0 t + C_2 \sin \omega_0 t + \frac{A}{2\omega_0} t \sin \omega_0 t.$$

Initial conditions $x(0) = x'(0) = 0 \Rightarrow$

$$x(t) = \frac{A}{2\omega_0} t \sin \omega_0 t.$$

Example: $\omega_0 = 5$, and $A = 2\omega_0 = 10$.

$$x(t) = t \sin 5t.$$

Solution

$$x(t) = \frac{A}{2\omega_0} t \sin \omega_0 t.$$

- Oscillation with increasing amplitude.
- First example of **resonance**.
 - ◇ Driving at the natural frequency can cause oscillations that grow out of control.

Forced, Damped Harmonic Motion

$$x'' + 2cx' + \omega_0^2 x = A \cos \omega t$$

- Assume the underdamped case, where $c < \omega_0$.

Homogeneous equation

$$x'' + 2cx' + \omega_0^2 x = 0$$

- Characteristic polynomial:

$$P(\lambda) = \lambda^2 + 2c\lambda + \omega_0^2$$

- Roots $\lambda = -c \pm \sqrt{c^2 - \omega_0^2} = -c \pm i\eta$ where $\eta = \sqrt{\omega_0^2 - c^2}$.

- Fundamental set of solutions

$$x_1(t) = e^{-ct} \cos \eta t \text{ and } x_2(t) = e^{-ct} \sin \eta t$$

Inhomogeneous equation

$$x'' + 2cx' + \omega_0^2 x = A \cos \omega t$$

- Use the complex method to find a particular solution. Solve

$$z'' + 2cz' + \omega_0^2 z = Ae^{i\omega t}.$$

Try $z(t) = ae^{i\omega t}$.

$$\begin{aligned} z'' + 2cz' + \omega_0^2 z &= [(i\omega)^2 + 2c(i\omega) + \omega_0^2]ae^{i\omega t} \\ &= P(i\omega)z \end{aligned}$$

Particular Solution

- Solution if $P(i\omega)z = Ae^{i\omega t}$.
- Solution is

$$\begin{aligned}z(t) &= \frac{1}{P(i\omega)} Ae^{i\omega t} \\ &= H(i\omega) Ae^{i\omega t}.\end{aligned}$$

- $H(i\omega) = \frac{1}{P(i\omega)}$ is called the **transfer function**.

$$\begin{aligned}P(i\omega) &= (i\omega)^2 + 2c(i\omega) + \omega_0^2 \\&= [\omega_0^2 - \omega^2] + 2ic\omega \\&= Re^{i\phi} \\&= R \cos \phi + iR \sin \phi.\end{aligned}$$

We need $R \cos \phi = \omega_0^2 - \omega^2$ and $R \sin \phi = 2c\omega$.

Thus

$$\begin{aligned}R &= \sqrt{(\omega_0^2 - \omega^2)^2 + 4c^2\omega^2} \\ \phi &= \operatorname{arccot} \left(\frac{\omega_0^2 - \omega^2}{2c\omega} \right).\end{aligned}$$

Transfer Function

$$\begin{aligned}H(i\omega) &= \frac{1}{P(i\omega)} \\ &= \frac{1}{R} e^{-i\phi} \\ &= G(\omega) e^{-i\phi}.\end{aligned}$$

- $G(\omega) = \frac{1}{R}$ is called the **gain**.

$$G(\omega) = \frac{1}{\sqrt{(\omega_0^2 - \omega^2)^2 + 4c^2\omega^2}}.$$

- ϕ is the **phase shift**.

Particular Solution

The complex particular solution is

$$\begin{aligned}z(t) &= H(i\omega)Ae^{i\omega t} \\ &= G(\omega)e^{-i\phi} \cdot Ae^{i\omega t} \\ &= G(\omega)Ae^{i(\omega t - \phi)}.\end{aligned}$$

The real particular solution is

$$\begin{aligned}x_p(t) &= \operatorname{Re}(z(t)) \\ &= G(\omega)A \cos(\omega t - \phi).\end{aligned}$$

General Solution

$$\begin{aligned}x(t) &= x_p(t) + x_h(t) \\ &= G(\omega)A \cos(\omega t - \phi) \\ &\quad + e^{-ct} [C_1 \cos \eta t + C_2 \sin \eta t].\end{aligned}$$

- Transient term.

- ◇ $x_h(t) = e^{-ct} [C_1 \cos \eta t + C_2 \sin \eta t]$.

- Steady-state solution.

- ◇ $x_p(t) = G(\omega)A \cos(\omega t - \phi)$.

Example

$\omega_0 = 1$, $c = 0.8$, $A = 1$, and $\omega = 3$.

$x(0) = 0$ and $x'(0) = 1$.

Steady-State Solution

$$x_p(t) = G(\omega)A \cos(\omega t - \phi).$$

- The forcing function is $A \cos \omega t$.
- The amplitude of the steady-state response is $G(\omega)$ times the amplitude of the forcing term.
- The steady-state response is oscillatory, but with a phase which is shifted by ϕ/ω .

Gain

$$G(\omega) = \frac{1}{\sqrt{(\omega_0^2 - \omega^2)^2 + 4c^2\omega^2}}$$

Set $\omega = s\omega_0$ and $c = D\omega_0/2$. Then

$$\omega_0^2 G(\omega) = \frac{1}{\sqrt{(1 - s^2)^2 + D^2 s^2}}$$

where

$$s = \frac{\omega}{\omega_0} \quad \text{and} \quad D = \frac{2c}{\omega_0}.$$